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It is known that the solution of the equation

$$\partial^2 \psi / \partial x^2 + \partial^2 \psi / \partial y^2 = F(\psi), \tag{1}$$

where the vorticity F is an arbitrary function of  $\psi$ , can be considered as an example of steady-state flow of an ideal fluid. If we suppose that the motion of an ideal incompressible fluid can be thought of as the threshold motion of a viscous fluid, the function  $F(\psi)$  in Eq. (1) can be replaced by a constant [1].

Let us consider the following simulation problem with cohesively selected piecewise-constant vorticity. In a bounded region D with boundary  $\Gamma$  it is necessary to find a continuously differentiable solution of the equation

$$\frac{\partial^2 \psi}{\partial x^2} + \frac{\partial^2 \psi}{\partial y^2} = \begin{cases} \omega, & \text{if } \psi < 0 \\ -\omega_1, & \text{if } \psi > 0 \end{cases}$$
 (2)

( $\omega$  and  $\omega_1$  are nonnegative constants) under the boundary condition

$$\psi|_{\Gamma} = \varphi(s). \tag{3}$$

If we set  $\omega_1 = 0$  in Eq. (2), we obtain an equation that describes the motion of an ideal fluid according to a previous scheme [2]. This type of flow for the case of a bounded region [3] and for the case of an unbounded region has been studied earlier [4-7].

The problem (2), (3) has the trivial solution

$$\psi = \varphi_0 + \frac{\omega_1}{2\pi} \int_{\mathcal{D}} G d\xi d\tau$$

where  $\varphi_0$  is a harmonic function satisfying condition (3) and G is Green's function of the region D of the Dirichlet problem for the Laplacian. In [3] it was proved that a nontrivial solution for the case  $\omega_1=0$  exists under particular conditions. We will derive a condition under which a nontrivial solution of the problem (2), (3) exists. A simpler bound than in [3] will be obtained from this condition for  $\omega_1=0$ .

Suppose  $\varphi$  (s)  $\leq$  C and let  $B_1$  be the circle of greatest radius, such that  $B_1 \subseteq D$  (without loss of generality we may assume that its center coincides with the coordinate origin), and let  $B_2$  be the circle of least radius with center at the origin, such that  $B_2 \supseteq D$ . The radius of  $B_1$  is  $R_1$  and that of  $B_2$ ,  $R_2$ . We have the following assertion: When

$$\omega - \frac{\omega_1 R_2^2}{R_1^2} e \geqslant \frac{4Ce}{R_1^2} \tag{4}$$

the problem (2), (3) has a nontrivial solution. Let us prove this assertion. If the circle  $B_1$  is taken as the region D and if we set  $\omega_1 = 0$  in Eq. (2), and let  $\varphi$  (s) = C +  $\omega_1 R_2^2/4$  in Eq. (3), whenever (4) holds, the problem has two nontrivial solutions (found explicitly). That is, in particular, there exists a circle  $B_a < B_1$  of radius a such that the corresponding solution is negative.

Let us consider the auxiliary problem

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$$\frac{\partial^2 \psi_n}{\partial x^2} + \frac{\partial^2 \psi_n}{\partial y^2} = \begin{cases} \omega, & \text{if} \quad x, y \in B_a \\ \frac{\omega}{2} (1 - \text{th } \psi_n n) - \frac{\omega_1}{2} (1 + \text{th } \psi_n n), & \text{if} \quad x, y \in D \setminus \overline{B}_a; \end{cases}$$
 (5)

$$\psi_n|_{\Gamma} = \varphi(s). \tag{6}$$

The solution will be found in the class of functions continuously differentiable in D. The problem (5), (6) is equivalent to the integral equation

$$\psi_n = \varphi_0 - \frac{\omega}{2\pi} \int_{B_a} \int G d\xi d\tau + \frac{1}{4\pi} \int_{D \setminus B_a} \left[ \omega_1 (1 + \operatorname{th} \psi_n n) - \omega (1 - \operatorname{th} \psi_n n) \right] G d\xi d\tau. \tag{7}$$

The Schauder theorem can be used to establish the existence of a solution of Eq. (7) for any n and  $x, y \in D \setminus \overline{B}_a$ . We substitute this solution in the right side of Eq. (7), thus defining the function  $\psi_n$  over all of D. The resulting function is the solution of the problem (5), (6). It follows from the properties of a potential-type integral that it has first derivatives in every fixed closed region  $\overline{B} \subset D$ ; these derivatives satisfy the Hölder condition, while the constant and exponent are independent of n.

We use the Arzela theorem to establish that the sequence  $\psi_n$  is compact in the space of continuously differentiable functions. Suppose the subsequence  $\psi_{nk}$  converges to a continuously differentiable function  $\psi^*$ . We will prove that  $\psi^*$  is a nontrivial solution of the problem (2), (3).

Suppose that  $y_0 \in D \backslash \overline{B}_a$  and  $\psi^*$  (x<sub>0</sub>, y<sub>0</sub>) > 0 at some point x<sub>0</sub>. It will then be greater than zero also in some circular neighborhood. We now consider Eq. (5) in this neighborhood and take its limit as  $n_k \to \infty$ , obtaining  $\partial^2 \psi^* / \partial x^2 + \partial^2 \psi^* / \partial y^2 = -\omega_1$ . It can be analogously proved that  $\partial^2 \psi^* / \partial x^2 + \partial^2 \psi^* / \partial y^2 = \omega$  at points at which  $\psi^* < 0$ . Further, when  $x,y \in B_a$ ,  $\partial^2 \psi^* / \partial x^2 + \partial^2 \psi^* / \partial y^2 = \omega$ . Let us prove that when x,  $y \in B_{a}$ ,  $\psi^* < 0$ . It follows from the properties of Green's function that

$$\frac{1}{2\pi} \int_{D} G d\xi d\tau \leqslant \frac{1}{2\pi} \int_{B_{1}} G_{B_{1}} d\xi d\tau \leqslant \frac{R_{2}^{2}}{4};$$

$$\int_{B_{2}} G d\xi d\tau \geqslant \int_{B_{2}} G_{B_{1}} d\xi d\tau (x, y \in B_{1}),$$
(8)

where  $G_{B_1}$  and  $G_{B_2}$  are Green's functions for the regions  $B_1$  and  $B_2$ , respectively. We find from Eqs. (7) and (8) that

$$\psi_n < V = C + \frac{\omega_1 R_2^2}{4} - \frac{\omega}{2\pi} \int_{B_a} G_{B_1} d\xi d\tau,$$

It follows from the definition of  $B_a$  that the function V is negative in  $B_a$ . Then  $\psi_n$ , that is,  $\psi^*$ , are both negative in  $B_a$ . The fact that  $\psi^*$  satisfies the equation as we pass through the boundary of  $B_a$  follows from its smoothness.

We set  $\omega_1 = 0$  in Eq. (4), obtaining the condition  $\omega \ge 4\text{Ce}/\text{R}_1^2$  under which there exists a nontrivial solution of the problem describing flow in the M. A. Lavrent'ev scheme for the case of a bounded region.

## LITERATURE CITED

- 1. G. K. Batchelor, Introduction to Fluid Dynamics, Cambridge University Press (1967).
- 2. M. A. Lavrent'ev, Variational Methods in Boundary-Value Problems for Systems of Elliptical-Type Equations [in Russian], Izd. Akad. Nauk SSSR, Moscow (1962).
- 3. M. A. Gol'dshtik, "Mathematical model of separated flows of an incompressible fluid," Dokl. Akad. Nauk SSSR, 147, No. 6 (1962).
- 4. A. B. Shabat, "Two attachment problems," Dokl. Akad. Nauk SSSR, 195, No. 6 (1963).
- 5. A. B. Shabat, "A scheme for plane motion of a fluid for a fracture on the floor," Zh. Prikl. Mekh. Tekh. Fiz., No. 4 (1962).
- 6. S. N. Antontsev and V. D. Lelyukh, "Certain attachment problems for eddy and potential flows," in: Continuum Dynamics [in Russian], Vol. 1, Izd. Inst. Gidrodinam., Sibirsk. Otd., Akad. Nauk SSSR, Novosibirsk (1969).
- 7. P. I. Plotnikov, "Solvability of a class of problems for attachment of eddy and potential flows," in: Continuum Dynamics [in Russian], Vol. 3, Izd. Inst. Gidrodinam., Sibirsk. Otd., Akad. Nauk SSSR, Novosibirsk (1969).